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Localized state and charge transfer in nitrogen-doped graphene

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Nitrogen-doped epitaxial graphene grown on SiC(000 $\bar{1}$) was prepared by exposing the surface to an atomic nitrogen flux. Using scanning tunneling microscopy and scanning tunneling spectroscopy (STS), supported by density functional theory (DFT) calculations, the simple substitution of carbon with nitrogen atoms has been identified as the most common doping configuration. High-resolution images reveal a reduction of local charge density on top of the nitrogen atoms, indicating a charge transfer to the neighboring carbon atoms. Local STS spectra clearly evidenced the energy levels associated with the chemical doping by nitrogen, localized in the conduction band. Various other nitrogen-related defects have been observed. The bias dependence of their topographic signatures demonstrates the presence of structural configurations more complex than substitution as well as hole doping.

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Graphene has been proposed as a promising alternative to silicon-based electronics for some applications. However, reliable control of its electronic properties, for example, by chemical doping, is still a challenging task.^{1–3} For carbon-based materials, the incorporation of nitrogen in the lattice is a natural choice because of its ability to form covalent bonds and to modify the electronic properties of sp^2 carbon locally, with minor structural perturbations.^{4–6} N-doped graphene also offers interesting prospects for various other applications, including biosensing,⁷ field emission,⁸ lithium incorporation,^{9,10} or transparent electrodes.¹¹

The substitution of some carbon atoms by nitrogen is expected to give rise to donor states and then to n -type doping.^{6,12} The synthesis of chemically modified graphene has been achieved either by direct growth of modified layers^{12–15} or by postgrowth treatment of pristine graphene.^{7,8,16–18} These studies revealed the presence of several atomic configurations for the nitrogen atoms: substitutional (“graphitic”), pyridine-like, or pyrrolelike N.⁸

However, a clear correlation between the synthesis methods and the atomic configuration of the chemically modified graphene, on the one hand, and between the atomic configuration and the electronic properties, on the other hand, remains a challenging task. A step in this direction has been achieved very recently by Zhao *et al.*¹³ through scanning tunneling microscopy (STM) and scanning tunneling spectroscopy (STS) investigations of nitrogen-doped graphene prepared by chemical vapor deposition on a copper substrate with NH_3 gas in the feedstock. These authors determined the atomic configuration of the nitrogen atoms to be predominantly (90%) a simple substitution (“graphitic” nitrogen), with the majority of dopants located on the same carbon sublattice of graphene. This work left open some key questions relative to the interpretation of the experimental data. For instance, simulated STM images exhibit a depletion above the N atom (see also Zheng *et al.*⁶) whereas experimental results have

no such central feature. More importantly, the N-induced donor energy level has not been evidenced in their STS measurements.

We present here a STM and STS study of N-doped graphene samples on SiC(000 $\bar{1}$) obtained by postsynthesis treatment. Our STM images provide clear evidence of the presence of substitutional nitrogen atoms together with more complex structures presenting well-defined topographical features. Local spectroscopy, supported by simulations based on density functional theory (DFT), reveals a localized donor state in the electronic density of states related to the substitutional nitrogen atoms. Based on STM images at different biases, we also observed and analyzed a change in the STM topographic image with respect to the tip-graphene distance that reconciles experiment and simulation. Doping configurations different from substitution are also analyzed.

Our graphene samples were prepared from C-terminated n -type 6H-SiC(000 $\bar{1}$) wafers, following procedures available in the literature^{19–23} which lead to multilayered graphene with misoriented, decoupled layers.²⁴ N doping was achieved by exposure to an atomic nitrogen flux produced by a remote radiofrequency plasma source fed with N_2 . The samples were then analyzed using a low-temperature STM (LT-STM) working under UHV conditions (see Supplemental Material and Ref. 11 therein²⁴). In our experimental setup, the plasma generator is not in the close vicinity of the sample, so that only N radicals (N^*) with thermal energy (and no accelerating voltage) interact with graphene. Thus, only the topmost layer is affected by the plasma treatment. This postsynthesis doping method has the advantage over direct growth of modified layers that a well-defined domain could be doped, the other part of the sample being kept pristine.

Figure 1 presents the graphene surface before (a) and after (b) N^* exposure. A defect-free honeycomb lattice with an

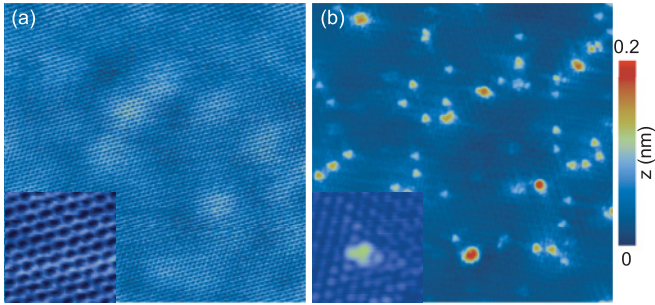


FIG. 1. (Color online) Comparison of STM images ($15 \times 15 \text{ nm}^2$) of graphene before (a) and after (b) N^* exposure. Inset of (a): Honeycomb lattice of pristine graphene. Inset of (b): $2 \times 2 \text{ nm}^2$ image of a nitrogen dopant. Tunneling parameters: (a) $V_s = -0.3 \text{ V}$, $I = 15 \text{ nA}$; (b) $V_s = -0.5 \text{ V}$, $I = 500 \text{ pA}$.

interatomic distance of 1.4 \AA is observed for the pristine graphene. After N treatment, many localized features appear as bright dots on the images, demonstrating the effect of the exposure to nitrogen radicals. Moreover, Auger electron spectroscopy (AES) confirms that the defects are related to nitrogen doping sites [no other chemical species (e.g., oxygen) are detected in AES] while Raman spectroscopy confirms the presence of defects sites after surface treatment (see Supplemental Material and Ref. 10 therein²⁴). By visual inspection of the STM images, a defect concentration ranging from 0.5% to 1.1% is deduced, depending on the exposure time (0.6% for the figures of the present paper).

A closer look at Fig. 1(b) reveals that different doping configurations are present. Approximately 75% of them display a triangular shape consisting of a bright spot (approximately 0.4 nm wide) with a threefold symmetry [inset of Fig. 1(b)]. These images are very similar to the one presented in Ref. 13 and are assigned to substitutional N atoms. No preferential orientation of the trigonal pattern is observed and, consequently, both sublattices of the graphene are affected by the N treatment. The graphene honeycomb pattern remains also unaltered outside of the vicinity of the defect, demonstrating a very local perturbation of the graphene layer. Other typical high-resolution images of this defect are shown in Figs. 2(a)–2(e) for biases ranging from $V_s = -0.4$ to $+0.5 \text{ V}$ (see Supplemental Material for a more complete overview²⁴). We also systematically observe that the corrugation at the doping site is more pronounced at positive biases, as compared to images recorded at negative biases. At this stage, the doping of graphene with N atoms can then be associated to spatially localized electronic states in the conduction band.

Ab initio calculations were performed to obtain the local electronic structure and the STM images using a Tersoff-Hamann approach. Calculations have been performed on a (9×9) supercell (0.6% of N atoms) and a (10×10) supercell (0.5% of N atoms) with a localized basis as implemented in the SIESTA package²⁵ (see more details in the Supplemental Material and Refs. 12–19 therein²⁴). The simulated patterns [Figs. 2(f) and 2(g); see also Refs. 6 and 13] present a dark (low) spot above the doping atom, surrounded by bright dots, corresponding to the adjacent C atoms. The central dark spot

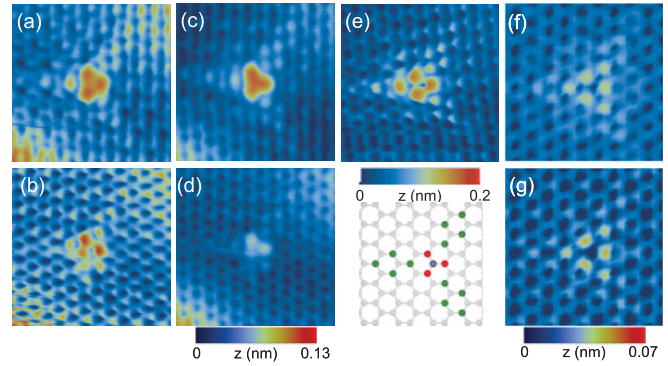


FIG. 2. (Color online) Topographic images ($2.5 \times 2.5 \text{ nm}^2$) of a substitutional nitrogen atom at (a) $V_s = +0.2 \text{ V}$, $I = 200 \text{ pA}$, (b) $V_s = -0.2 \text{ V}$, $I = 100 \text{ pA}$, (c) $V_s = +0.5 \text{ V}$, $I = 700 \text{ pA}$, (d) $V_s = -0.4 \text{ V}$, $I = 100 \text{ pA}$, and (e) $V_s = +0.35 \text{ V}$, $I = 800 \text{ pA}$. In the schematic view, the central small dot (blue online) corresponds to the nitrogen atom, and the darker and lighter dots (red and green online) correspond to the carbon atomic sites around which the density of states is higher with decreasing values from dark to light (red to green online), as suggested by the experimental images. (f) and (g) Simulations for N substitution in a 10×10 graphene supercell at (f) $V_s = +0.5 \text{ V}$ and (g) $V_s = -0.5 \text{ V}$.

has been explained by charge transfer from the N atom to the neighboring C atoms, which results in a smaller spatial extension of the electronic states associated with the N atom in the direction perpendicular to the layer compared to the one associated with the C atoms forming the C-N bonds.⁶ This central “hole” is not observed in most STM images (Fig. 2 and Ref. 13) but both the weak dependence of the pattern shape with the bias voltage and the more intense corrugation for positive biases are reproduced by simulations. Generally speaking, the limit of our computation procedure, besides the intrinsic approximations of DFT, are the simplified tip (a “metallic” s shape atomic orbital) and, more important in the present study, the small distance between the tip and the atomic layer. For numerical reasons, this distance is smaller than $3\text{--}4 \text{ \AA}$ in the simulations, i.e., close to the point contact and underestimated with respect to the experimental distance.

Interestingly, the hole at the center of the defect (the N atom in chemical substitution) appears in some of the experimental images, as in Fig. 2(b), i.e., at low negative bias (see also the Supplemental Material²⁴). By varying more extensively the experimental conditions, the same pattern has been observed for $V_s = +0.35 \text{ V}$ and high current [Fig. 2(e)]. In all these cases (low V_s or high current), the tip is close to the surface and a better agreement with the simulations is expected, and is indeed observed. The tip is also closer to the surface for negative bias than for positive bias because of the localization of the energy state associated with the defect in the conduction band (see below). This explains why a plain triangular pattern is observed for $V_s = +0.5 \text{ V}$ whereas a hollow pattern is observed at $V_s = -0.2 \text{ V}$.

Further insight regarding the local electronic structure of the dopant can be gained through STS spectra [Fig. 3(a), with the corresponding image on Fig. 3(b)]. The spectra display two minima: one at the Fermi energy (0 V) associated with phonon-mediated inelastic channels,^{13,26} and one corresponding to the

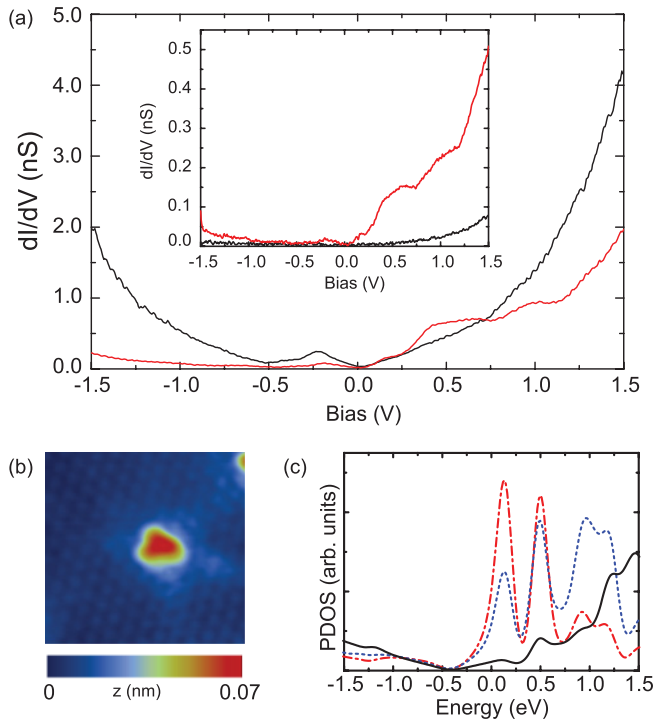


FIG. 3. (Color online) Comparison of scanning tunneling spectra between graphene (black curve) and the simple substitution (gray curve, red online). (a) Spectra taken with the feedback loop active when moving from one spot to another. Inset: Spectra taken with the feedback loop off when moving from one spot to another. (b) Topographic image of the defect on which the spectra on (a) have been taken. Tunneling parameters: $V_s = +1.0$ V, $I = 500$ pA. (c) Simulated partial DOS for a graphene 9×9 supercell (162 atoms) including a single N substitution. The PDOS far away from the nitrogen atom (C_{bulk}), on the neighboring C atoms (C_1) and on the nitrogen atom are represented by the solid, dashed, and dotted-dashed (black, blue, and red online) curves, respectively.

Dirac energy at -0.5 V. The relationship between the charge carrier concentration n , the Dirac energy E_D , and the Fermi velocity (v_F) $n = \frac{E_D^2}{\pi(\hbar v_F)^2}$ leads to $n = 18 \times 10^{12}$ electrons per cm^2 for $v_F = 10^6$ m/s. Assuming 0.6% of nitrogen atoms, the charge transfer can be estimated to 0.8 electron per dopant atom. From a theoretical viewpoint, our simulations for a 0.6% doping [Fig. 3(c)] exhibit E_D at -0.42 eV and then a smaller charge transfer (0.55 electron per N atoms). These values are larger than those reported by Zhao *et al.*¹³ and the discrepancy can be explained by a number of potential sources of uncertainties: the accuracy in the determination of the nitrogen concentration as well as the position of the Dirac point, the presence of several types of doping sites, and the uncertainty on the value of the Fermi velocity^{27,28} on the experimental side or the absence of quasiparticle corrections,²⁹ and the regular spacing between the N sites in the supercell technique for the DFT calculations.

The comparison between the dI/dV spectra at the nitrogen doping sites and far from it [Fig. 3(a)] shows that a broad peak centered around $+0.5$ eV appears in the vicinity of the nitrogen atom [dI/dV spectra give a quantity proportional to the local density of states (LDOS)]. This is an experimental

determination of the energy level of the localized state in N-doped graphene. The simulated partial density of states (PDOS) located on the N atom, and on a C atom close to and far from the nitrogen (C_1 and C_{bulk} atoms, respectively) are displayed in Fig. 3(c). The states related to the graphitic nitrogen are clearly obtained in the conduction bands but display a double-peak structure at 0.15 and 0.50 eV. The nitrogen-nitrogen interactions in the periodical structure used in the calculation could lead to a splitting of the donor state that would not be present for randomly distributed defects.³⁰ The absence of quasiparticle corrections²⁹ or the reduced tunneling current for low bias²⁶ may also explain why the low-energy states are not observed. Besides, experimental STS spectra obtained for a tip above the center of the defect probably also probe the PDOS of the C_1 atoms. As the N PDOS, the C_1 PDOS presents a double peak at 0.15 and 0.50 eV but with a larger amplitude of the second one. Moreover, a state localized at the C_1 atoms appears around 1 eV and can be related to a second feature in the experimental STS.

Coming back to the experimental STS spectra, at negative bias, the dI/dV signal is found to be lower above the nitrogen atom than above the graphene, as shown in Fig. 3(a). At positive bias both STS spectra have comparable intensities. To reconcile those measurements with the higher corrugation above the dopant site (Figs. 1 and 2), we have to keep in mind that spectra are measured with initial conditions corresponding to the setpoint used for the STM image. As a consequence, the tip is located at a larger distance from the atomic plane when a spectrum is measured above the dopant site, compared with a STS measurement on the graphene layer and the two spectra cannot be quantitatively compared. As it is expected that the nitrogen atom lies in the plane of the graphene sheet, as demonstrated by DFT simulations performed on single⁶ and bilayer N-doped graphene,³¹ the observed corrugation is a purely electronic effect. In such cases, the intensity of the dI/dV spectra is artificially reduced above the N atom due to the higher tip position. In order to overcome this difficulty of interpretation, we have measured some spectra while scanning with the STM feedback loop off, i.e., at a constant tip height. In these conditions, STS spectra exhibit a very different behavior [inset in Fig. 3(a)]. For negative bias, the spectrum recorded above the dopant is slightly larger than above the graphene layer, in agreement with what is observed on the images. For positive bias, the difference is even more spectacular with a LDOS measured at the doping site up to five times more intense than the one of graphene, strongly confirming that the doping states lie in the conduction band.

Besides the substitutional doping, other types of defects have been frequently observed, examples of which are displayed in Fig. 4. Some defects appear higher at negative bias [Figs. 4(a) and 4(c)] than at positive bias [Figs. 4(b) and 4(d)], contrary to what was observed for substitution. Interestingly, a simple substitutional N atom is also observed on top of Figs. 4(a)–4(d) close to the complex defects. These configurations allow a direct comparison of the bias dependence of the simple substitutional N (higher at positive bias) and of the other defects (higher at negative bias). These unique dopant sites can then be associated with hole doping and with a localized state lying in the valence band.^{4,6} Furthermore, the experimental signature (a triangular symmetry with extended oscillations of

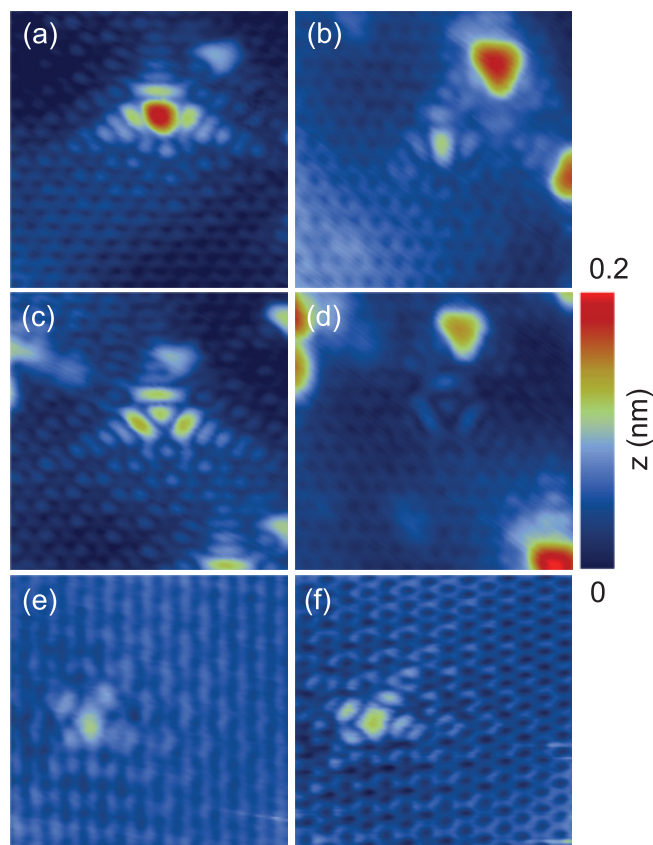


FIG. 4. (Color online) STM images of complex doping sites ($3 \times 3 \text{ nm}^2$). (a), (c), and (e) correspond to three different defects measured at a bias voltage of -0.5 V . (b), (d), and (f) correspond respectively to the same defects measured at $+0.5 \text{ V}$.

the wave function) is consistent with the behavior predicted for a configuration involving three N atoms forming a pyridinic configuration⁶ or with a single vacancy.³² Other defects,

such as the one reported in Figs. 4(e) and 4(f), present a pattern with a rectangular symmetry and exhibit an *n*-type characteristic (higher corrugation at positive bias). However, further investigations based on STS analysis are necessary to reach a definitive conclusion on the atomic configurations of these complex doping sites.

In summary, the N doping of graphene was achieved by exposing epitaxial samples to an atomic nitrogen flux. The most common ($\sim 75\%$) doping configuration is found to be substitutional (graphitic) N atoms. Local STS spectra have provided experimental evidence of a localized donor state, a specific signature of N substitution. STM topographical images revealed a depletion above the N sites, characteristic of a charge transfer between the N atom and the adjacent neighboring C atoms. Finally, other N doping sites with a more complex atomic structure than a simple substitution were observed, exhibiting an acceptorlike character. The presented postsynthesis treatment opens the door to local tuning of electronic properties of graphene, which is a prerequisite for the development of carbon-based electronics.

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¹A. H. Castro Neto, F. Guinea, N. M. R. Peres, K. S. Novoselov, and A. K. Geim, *Rev. Mod. Phys.* **81**, 109 (2009).

²H. Liu, Y. Liu, and D. Zhu, *J. Mater. Chem.* **21**, 3335 (2011).

³M. S. Fuhrer, C. N. Lau, and A. H. MacDonald, *MRS Bull.* **35**, 289 (2010).

⁴C. P. Ewels and M. Glerup, *J. Nanosci. Nanotechnol.* **5**, 1345 (2005).

⁵A. Lherbier, X. Blase, Y.-M. Niquet, F. Triozon, and S. Roche, *Phys. Rev. Lett.* **101**, 036808 (2008).

⁶B. Zheng, P. Hermet, and L. Henrard, *ACS Nano* **7**, 4165 (2010).

⁷Y. Wang, Y. Shao, D. Watson, J. Li, and Y. Lin, *ACS Nano* **4**, 1790 (2010).

⁸N. Soin, S. S. Roy, S. Roy, K. S. Hazra, D. S. Misra, T. H. Lim, C. J. Hetherington, and J. A. McLaughlin, *J. Phys. Chem. C* **115**, 5366 (2011).

⁹A. L. M. Reddy, A. Srivastava, S. R. Gowda, H. Gullapalli, M. Dubey, and P. M. Ajayan, *ACS Nano* **4**, 6337 (2010).

¹⁰Y. J. Cho, H. S. Kim, H. Im, Y. Myung, G. B. Jung, C. Lee, J. Park, M.-H. Park, J. Cho, and H. S. Kang, *J. Phys. Chem. C* **115**, 9451 (2011).

¹¹A. Kasry, M. A. Kuroda, G. Martyna, G. Tulevski, and A. a. Bol, *ACS Nano* **4**, 3839 (2010).

¹²D. Wei, Y. Liu, Y. Wang, H. Zhang, L. Huang, and G. Yu, *Nano Lett.* **9**, 1752 (2009).

¹³L. Zhao, R. He, K. T. Rim, T. Schiros, K. S. Kim, H. Zhou, C. Gutiérrez, S. P. Chockalingam, C. J. Arguello, L. Pálová *et al.*, *Science* **333**, 999 (2011).

¹⁴L. S. Panchakarla, K. S. Subrahmanyam, S. K. Saha, A. Govindaraj, H. R. Krishnamurthy, U. V. Waghmare, and C. N. R. Rao, *Adv. Mater.* **26**, 4726 (2009).

- ¹⁵J. C. Meyer, S. Kurasch, H. J. Park, V. Skakalova, D. Künzel, A. Groß, A. Chuvilin, G. Algara-Siller, S. Roth, T. Iwasaki *et al.*, *Nat. Mater.* **10**, 209 (2011).
- ¹⁶B. Guo, Q. Liu, E. Chen, H. Zhu, L. Fang, and J. R. Gong, *Nano Lett.* **10**, 4975 (2010).
- ¹⁷Y.-C. Lin, C.-Y. Lin, and P.-W. Chiu, *Appl. Phys. Lett.* **96**, 133110 (2010).
- ¹⁸H. M. Jeong, J. W. Lee, W. H. Shin, Y. J. Choi, H. J. Shin, J. K. Kang, and J. W. Choi, *Nano Lett.* **11**, 2472 (2011).
- ¹⁹I. Forbeaux, J.-M. Themlin, and J.-M. Debever, *Surf. Sci.* **442**, 9 (1999).
- ²⁰W. A. de Heer, C. Berger, X. Wua, P. N. First, E. H. Conrada, X. Li, T. Li, M. Sprinklea, J. Hassa, M. L. Sadowski *et al.*, *Solid State Commun.* **143**, 92 (2007).
- ²¹F. Varchon, P. Mallet, L. Magaud, and J.-Y. Veuillen, *Phys. Rev. B* **77**, 165415 (2008).
- ²²F. Hiebel, P. Mallet, F. Varchon, L. Magaud, and J.-Y. Veuillen, *Phys. Rev. B* **78**, 153412 (2008).
- ²³U. Starke and C. Riedl, *J. Phys.: Condens. Matter* **21**, 134016 (2009).
- ²⁴See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevB.85.161408> and Refs. 1–9 therein for details on the sample preparation and on the sample properties.
- ²⁵D. Sánchez-Portal, P. Ordejón, E. Artacho, and J. M. Soler, *Int. J. Quantum Chem.* **65**, 453 (1997).
- ²⁶Y. Zhang, V. W. Brar, F. Wang, C. Girit, Y. Yayon, M. Panlasigui, A. Zettl, and M. F. Crommie, *Nat. Phys.* **4**, 627 (2008).
- ²⁷J. Hicks, M. Sprinkle, K. Shepperd, F. Wang, A. Tejada, A. Taleb-Ibrahimi, F. Bertran, P. Le Fèvre, W. A. de Heer, C. Berger *et al.*, *Phys. Rev. B* **83**, 205403 (2011).
- ²⁸D. L. Miller, K. D. Kubista, G. M. Rutter, M. Ruan, W. A. de Heer, P. N. First, and J. A. Stroscio, *Science* **324**, 924 (2009).
- ²⁹Quasiparticle corrections have been demonstrated to lead to a 15% correction of the Fermi velocity of graphene (Ref. 33). Even if the *GW* calculations have not been performed here due to the large size of the supercell, the correction is expected to be small.
- ³⁰P. Lambin, H. Amara, F. Ducastelle, and L. Henrard (unpublished).
- ³¹S.-O. Guillaume, B. Zheng, J.-C. Charlier, and L. Henrard, *Phys. Rev. B* **85**, 035444 (2012).
- ³²H. Amara, S. Latil, V. Meunier, P. Lambin, and J.-C. Charlier, *Phys. Rev. B* **76**, 115423 (2007).
- ³³P. E. Trevisanutto, C. Giorgetti, L. Reining, M. Ladisa, and V. Olevano, *Phys. Rev. Lett.* **101**, 226405 (2008).