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Published in:

Nuclear Instruments and Methods in Physics Research, Section B: Beam Interactions with Materials and Atoms

DOI:

[10.1016/j.nimb.2005.06.140](https://doi.org/10.1016/j.nimb.2005.06.140)

Publication date:

2005

Document Version

Early version, also known as pre-print

[Link to publication](#)

Citation for published version (HARVARD):

Colaux, JL & Terwagne, G 2005, 'Simultaneous depth profiling of the ^{12}C and ^{13}C elements in different samples using (d,p) reactions', *Nuclear Instruments and Methods in Physics Research, Section B: Beam Interactions with Materials and Atoms*, vol. 240, no. 1-2, pp. 429-433. <https://doi.org/10.1016/j.nimb.2005.06.140>

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Simultaneous depth profiling of the ^{12}C and ^{13}C elements in different samples using (d,p) reactions

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Available online 2 August 2005

Abstract

Nuclear reactions (d,p) are often used to perform depth profiling of light elements in solids. In particular, protons coming from $^{12}\text{C}(\text{d},\text{p}_0)^{13}\text{C}$ and $^{13}\text{C}(\text{d},\text{p}_0)^{14}\text{C}$ reactions are emitted at very different energies. Consequently these two reactions can be used to depth profile ^{12}C and ^{13}C simultaneously. Nevertheless the cross-section of $^{13}\text{C}(\text{d},\text{p}_0)^{14}\text{C}$ reaction is 10 times smaller than the $^{12}\text{C}(\text{d},\text{p}_0)^{13}\text{C}$ one. So, the geometry of detection must be judiciously chosen in order to depth profile these two elements with a high sensitivity and good resolution.

In the framework of this study we have performed 400 keV ^{13}C ions implantation into polished copper substrates at different temperatures and implanted doses with a 2 MV Tandem accelerator. Using the reactions described above, we have studied the evolution of ^{13}C depth profile as a function of implanted doses and temperature. We have also determined the origin of surface contamination that appears during the implantation process.

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PACS: 25.45.De; 29.30.Ep

Keywords: ^{12}C ; ^{13}C ; (d,p); Profile concentration; NRA

1. Introduction

Since 1985 and the fullerene (C_{60}) discovery [1], there has been an intensive effort driven by the scientific community to study the so-called novel car-

bon nanostructures (fullerenes, nanotubes, carbon onions). At present there are numerous techniques allowing to synthesize carbon onions [2–5]. One of them developed by Cabioch et al. in the early 1990s uses high dose carbon ion implantation into hot substrates (like copper or silver) [4,6–9]. Recently the same team has been shown by successive ^{13}C and ^{12}C low energy implantation into silver that depth distribution of implanted carbon has

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a Gaussian-like shape and there is no preferential carbon diffusion towards the surface during the implantation process [10].

In this paper, we study high energy implantation of ^{13}C into copper and we try to explain the depth distributions observed as a function of implanted dose and temperature. On the other hand, it is well known that build-up of carbon occurs during ion implantation. This phenomenon is directly related to residual vacuum and implantation parameters and we try to show that it is responsible to the surface contamination observed on the implanted samples. We optimized the detector geometry in order to depth profile ^{12}C and ^{13}C with high sensitivity and good resolution by using nuclear reactions induced by deuteron.

2. Experimental details

Polished copper substrates were fixed into a furnace mounted in an implantation chamber to perform different $^{13}\text{C}^+$ ions implantations at different temperatures (200, 370 and 600 °C). Each one was performed in residual pressure less than 7×10^{-4} Pa. $^{13}\text{C}^+$ ions of 400 keV were produced by ALTAIS¹ the 2 MV Tandem accelerator installed at LARN (Namur). About 4 h of implantation was necessary to obtain an average dose of 5×10^{17} at. cm⁻² on a rectangular surface (5 × 6 mm²) with ion beam intensity of about 1.6 μA. Due to the fact that the beam was not homogeneous, we are able to measure the areal density at different places on the implanted specimens using the reactions described below.

In order to characterize the implanted samples in one measurement, we used non-resonant nuclear reactions (NRA) induced by 1.050 MeV deuteron beam (1 × 1 mm²). Two silicon surface barrier detectors were mounted at 135° and 165° relative to the incident beam. The NRA detector at 135° measured protons emitted by $^{12}\text{C}(\text{d},\text{p}_0)^{13}\text{C}$ and $^{13}\text{C}(\text{d},\text{p}_0)^{14}\text{C}$ reactions. A 12 μm MylarTM absorber foil was placed in front of the detector to

stop the intense flow of backscattered ions. Since the cross-section of the $^{13}\text{C}(\text{d},\text{p}_0)^{14}\text{C}$ is 10 times lower (about 3 mb sr⁻¹) [11] than the cross-section of $^{12}\text{C}(\text{d},\text{p}_0)^{13}\text{C}$ [12–14], we have chosen to use a large solid angle (25.7 msr) in order to obtain a good statistic while limiting the acquisition time to about 10 min. The RBS detector placed at 165° was collimated with a very small solid angle (0.18 msr) in order to detect backscattered deuterons flow without any absorber. The backscattered deuterons on copper were used as a monitor to measure incident particles [15]. Each of the three samples was measured in five different places in order to evaluate the non-uniformity of carbon beam irradiation. The acquisition time for each measurement was monitored by integrating the current beam to obtain an accumulated charge of 15 μC.

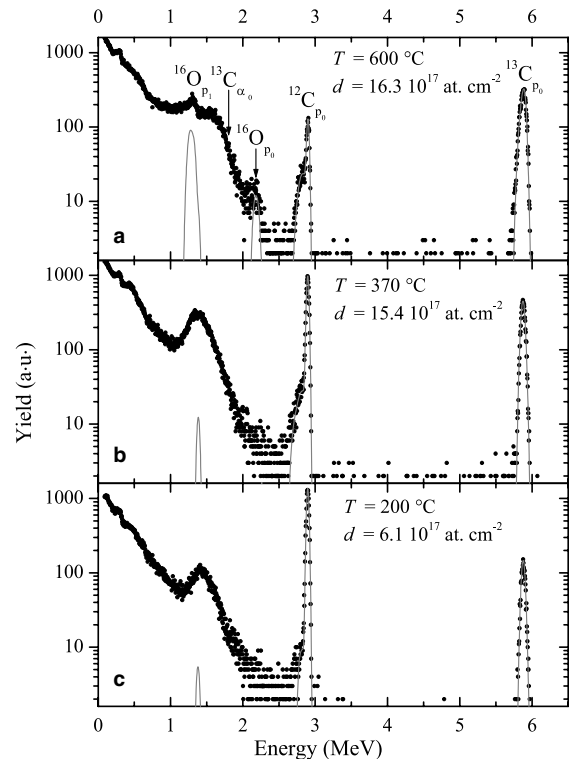


Fig. 1. Experimental NRA spectra of copper samples implanted with $^{13}\text{C}^+$ at various temperatures (T) and fluences (d) recorded in the NRA detector. Symbols represent the experimental spectra and the black lines represent simulations realized with SIMNRA code.

¹ Accélérateur Linéaire Tandem pour l'Analyse et l'Implantation des Solides.

3. Results

The measurements were realized at different places on each sample and show that the implanted dose varied in a wide range (from 2×10^{17} to 16.3×10^{17} at. cm⁻²) according to the position on the irradiated zone. These doses were calculated by simulation of the experimental spectra with the SIMNRA code [16]. As indicated in the previous section, the average implanted dose is 5×10^{17} ¹³C cm⁻² and we have reactions with a good yield to reveal the non-uniformity which is not the first objective of this paper.

Fig. 1 shows typical experimental NRA spectra of copper samples implanted with ¹³C at various temperatures recorded in the NRA detector. The ¹²C_{p0} at 2.9 MeV and the ¹³C_{p0} at 5.9 MeV are due to nuclear reactions ¹²C(d,p₀)¹³C and ¹³C(d,p₀)¹⁴C respectively. We have adopted the same identification for the reactions induced

by deuterons (i.e. ¹⁶O_{p0} as nuclear reaction ¹⁶O(d,p₀)¹⁷O). We can observe that the intensity of the ¹³C_{p0} peak is proportional to the implanted dose. The shape of this peak is wider for the sample implanted at high temperature (Fig. 1(a)). This effect due to diffusion phenomenon during implantation will be discussed later.

The intensity of the ¹²C_{p0} peak reveals the build-up phenomenon during the implantation. At temperature lower than 370 °C (Fig. 1(b) and (c)) we can observe a very intense contribution of ¹²C at the surface while ¹²C contamination is strongly reduce when the implantation temperature is 600 °C (Fig. 1(a)). As it can be seen in Fig. 1, a shoulder at low energy of the ¹²C_{p0} peak reveals the high sensitivity of this reaction. We can measure the areal density of ¹²C in the bulk of the implanted samples.

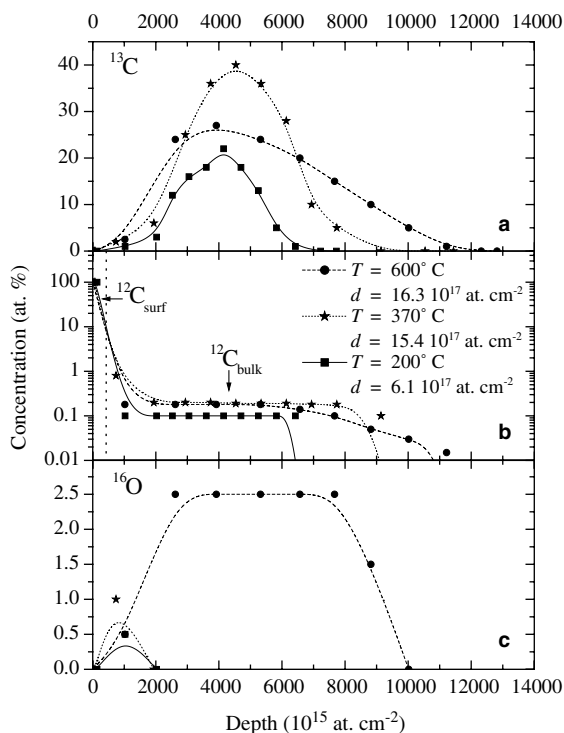


Fig. 2. Depth distributions of (a) ¹³C, (b) ¹²C and (c) ¹⁶O for 400 keV ¹³C⁺ implantations into copper performed at different temperatures (*T*) and for various fluences (*d*).

Table 1

Equivalent thickness (expressed in 10¹⁵ at. cm⁻²) of different components of implanted samples according to SIMNRA simulations

	Sample implanted at 200 °C	Sample implanted at 370 °C	Sample implanted at 600 °C
Implanted dose (¹³ C) (10 ¹⁵ at. cm ⁻²)	6.1×10^2	15.4×10^2	16.3×10^2
¹² C _{surf}	2.75×10^2	1.38×10^2	0.18×10^2

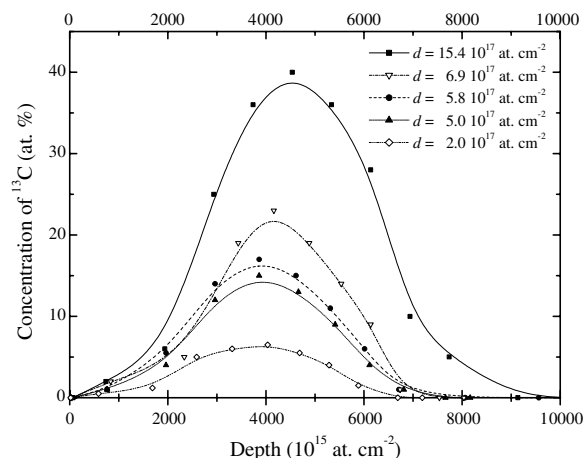


Fig. 3. Depth distributions of ¹³C for 400 keV ion implantation into copper performed at 370 °C with various fluences (*d*).

Table 2
Comparison between $(R_p)_{\text{exp}}$ and $(R_p)_{\text{th}}$ as function of implanted dose

Dose (10^{17} at. cm^{-2})	$[^{13}\text{C}]^{\text{max}}$ (at.%)	$(R_p)_{\text{exp}}$ (10^{15} at. cm^{-2})	Composition of target	$(R_p)_{\text{th}}$ (10^{15} at. cm^{-2})
15.4	40	4300	$\text{C}_{0.200}\text{Cu}_{0.800}$	4140
6.8	23	4040	$\text{C}_{0.115}\text{Cu}_{0.885}$	3930
5.8	17	4020	$\text{C}_{0.085}\text{Cu}_{0.915}$	3950
5.0	15	3960	$\text{C}_{0.075}\text{Cu}_{0.925}$	3880
2.0	6.5	3940	$\text{C}_{0.032}\text{Cu}_{0.968}$	3810

The fourth column indicates the composition of the target used to calculate the $(R_p)_{\text{th}}$ with the SRIM code.

At low energy (<2.5 MeV) the spectra for the high temperature (600 °C) and the low temperatures (200 °C and 370 °C) are completely different. We can observe in Fig. 1(a) the peaks $^{16}\text{O}_{\text{p}_0}$ and $^{16}\text{O}_{\text{p}_1}$, due to the presence of oxygen, at 2.2 and 1.35 MeV respectively. The sensitivity of these reactions is limited by an interference with the $^{13}\text{C}(\text{d},\alpha_0)^{11}\text{B}$. In order to simulate the spectra at low energies, the cross-section of this reaction already measured by Marion et al. [17] for energies higher than 1 MeV should be measured at lower energies.

4. Discussion

Fig. 2 shows the depth distributions of ^{12}C , ^{13}C and ^{16}O . The concentration of samples has been sliced into layers containing all concerned elements, which are represented by the different points. The surface contribution ($^{12}\text{C}_{\text{surf}}$) is clearly observed in Fig. 2(b) and the areal density of ^{12}C has been calculated and reported in Table 1. This surface contribution is due to build-up of ^{12}C during the implantation. The sample implanted at 600 °C contains the lowest ^{12}C contamination. This effect is explained by the dependence of the sputtering yield to the temperature. As the implantation temperature is increasing, the sputtering yield is also increasing [18,19].

At medium implantation temperature (370 °C) the ^{13}C depth profiles show a Gaussian shape (Fig. 3). The projected range increases with the implanted dose. The experimental and theoretical projected ranges, calculated with SRIM2003 code [20], are reported in Table 2. The theoretical R_p values have been calculated for a target composition corresponding to the half maximum of carbon

contents (fourth column in Table 2). It is clearly seen in Table 2 that the agreement between experimental and calculated values is very good.

The ^{13}C depth profile at high temperature (600 °C) shows a dissymmetrical distribution (Fig. 2(a)) which is correlated to the ^{12}C depth distribution (Fig. 2(b)). The wide distribution of ^{13}C is explained by diffusion process during implantation. This effect should be confirmed with a theoretical model like TRIDYN code [21] adapted to take into account the diffusion process. The correlation between ^{13}C and $^{12}\text{C}_{\text{bulk}}$ distribution is also explained by diffusion of ^{12}C during the implantation process.

Moreover, the oxygen depth distribution of the high implantation temperature is correlated to ^{13}C depth distribution. This effect should also be confirm when $^{13}\text{C}(\text{d},\alpha_0)^{11}\text{B}$ will be measured at low energies.

5. Conclusion

We have developed an ideal geometry of detection in order to analyse simultaneously ^{12}C , ^{13}C and ^{16}O by using nuclear reactions induced by deuteron. Two detectors were used to perform these measurements. The first detector positioned in RBS geometry was used as a monitor and the second one placed in NRA geometry showed a large solid angle to depth profile ^{12}C , ^{13}C and ^{16}O simultaneously. Depth profiling of ^{12}C reveals a surface contamination due to ^{12}C build-up during implantation process and small bulk component which is correlated with ^{13}C depth distribution. The sensitivity of the $^{12}\text{C}(\text{d},\text{p}_0)^{13}\text{C}$ nuclear reaction allows to detect 0.1 at.% of ^{12}C . At low implantation temperature (<370 °C), the ^{13}C depth profiles show a

Gaussian shape with $(R_P)_{\text{exp}}$ close to theoretical range $(R_P)_{\text{th}}$ calculated with SRIM code. For high temperature of implantation (600 °C), the ^{13}C depth profile shows a diffusion process.

Finally we were able to make obvious that only first 1000×10^{15} at. cm^{-2} samples are weakly oxidized as long as implantation temperature is maintained below 370 °C. On the other hand oxidation is 10 times higher for implantation realized at 600 °C. Moreover ^{16}O depth profile seems to be correlated to ^{13}C depth distribution. In order to measure precisely oxygen, the cross-section of $^{13}\text{C}(d,\alpha_0)^{11}\text{B}$ nuclear reaction should be measured for projectiles energies below 1 MeV.

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